

Exclusive π^0 production at EIC of China within handbag approach*

S.V. Goloskokov^{1†} Ya-Ping Xie(谢亚平)^{2,3‡} Xurong Chen(陈旭荣)^{2,3,4§}

¹Laboratory of Theoretical Physics, Joint Institute for Nuclear Research, Dubna 141980, Moscow, Russia

²Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, China

³University of Chinese Academy of Sciences, Beijing 100049, China

⁴Institute of Quantum Matter, South China Normal University, Guangzhou 510006, China

Abstract: Exclusive π^0 electroproduction is analyzed within the handbag approach based on Generalized Parton Distribution (GPDs) factorization. We consider the leading-twist contribution together with the transversity effects. It is shown that the transversity GPDs H_T and \bar{E}_T are essential in the description of the π^0 cross section. Predictions for the future Electron-Ion Collider of China (EicC) energy range are provided. It is found that transversity dominance $\sigma_T \gg \sigma_L$, observed at low energies, is valid up to the EicC energy range.

Keywords: general parton distributions, nucleon structure, deep virtual meson production

DOI: 10.1088/1674-1137/ac878c

I. INTRODUCTION

The study of the hadron structure is one of the key problems in modern physics. Some time ago, while analyzing exclusive processes, a new object was proposed, Generalized Parton Distributions (GPDs) [1–3]. It was found that the exclusive processes at a large photon virtuality Q^2 , such as the deeply virtual Compton scattering (DVCS) [4–6] and deeply virtual meson production (DVMP) [7–9], factorize into hard subprocesses that can be calculated perturbatively and GPDs [4–6]. Generally, this factorization was proved in the leading-twist amplitude with a longitudinally polarized photon.

GPDs are complicated nonperturbative objects that depend on 3 variables, i.e., x - the momentum fraction of a proton carried by the parton, ξ - skewness, and t - momentum transfer. GPDs contain information about the longitudinal and transverse distributions of the partons inside the hadron. They give information on its 3D structure; see, e.g., [10].

In the forward limit ($\xi = 0, t = 0$), GPDs become equal to the corresponding parton distribution functions (PDFs). The form factors of hadrons can be calculated from GPDs through integration over x [4]. Using Ji sum rules [4], the parton angular momentum can be extracted. More information on GPDs can be found e.g. in [7, 11,

12].

The study of exclusive meson electroproduction is one of the most effective ways of accessing GPDs. An experimental study of π^0 production was performed by CLAS [13] and COMPASS [14]. These experimental datasets can be adopted to constrain the models of GPDs. Electron-Ion Colliders (EICs) are the next generation of colliders for investigating the nucleon structure in the future. USA and China both plan to build EICs in the next 20 years [15–17]. The GPDs are one of the most important aspects that need to be studied for the EICs [18].

The theoretical study of DVMP in terms of GPDs is often based on the handbag approach, in which, as mentioned before, the amplitudes factorize into the hard subprocess and GPDs [2–5]; see Fig. 1. This amplitude contains another non-perturbative object distribution amplitude, which can be used to probe the two-quark components of the meson wave function. One of the popular ways of constructing GPDs is using the so-called Double Distribution (DD) [19], which constructs ξ dependencies of GPDs and connects them with PDFs, modified by a t -dependent term. The handbag approach with the DD form of GPDs was successfully applied to light vector meson (VM) leptonproduction at high photon virtualities Q^2 [20–22] and pseudoscalar meson (PM) leptonproduction [23]. In this work, we compute the π^0 production by ap-

Received 15 June 2022; Accepted 8 August 2022; Published online 19 September 2022

* Supported by the Chinese Academy of Sciences president's international fellowship initiative (2021VMA0005) and Strategic Priority Research Program of Chinese Academy of Sciences (XDB34030301)

[†] E-mail: goloskkv@theor.jinr.ru

[‡] E-mail: xieyaping@impcas.ac.cn, Corresponding author

[§] E-mail: xchen@impcas.ac.cn



Content from this work may be used under the terms of the Creative Commons Attribution 3.0 licence. Any further distribution of this work must maintain attribution to the author(s) and the title of the work, journal citation and DOI. Article funded by SCOAP³ and published under licence by Chinese Physical Society and the Institute of High Energy Physics of the Chinese Academy of Sciences and the Institute of Modern Physics of the Chinese Academy of Sciences and IOP Publishing Ltd

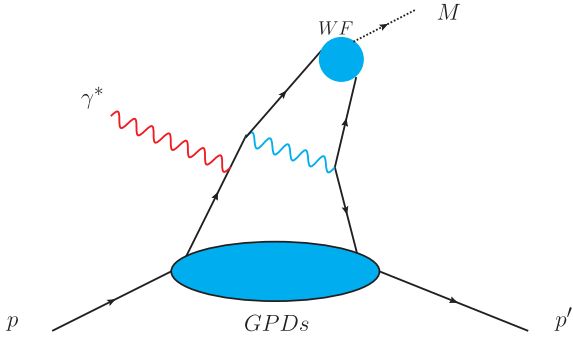


Fig. 1. (color online) The handbag diagram for the meson electroproduction of protons.

plying the handbag approach to the kinematics for EIC in China (EicC). Our prediction for π^0 production is helpful for estimating the meson cross section at EicC in the future.

In the leading twist approximation, the amplitudes of the pseudoscalar meson leptonproduction are sensitive to the GPDs \tilde{H} and \tilde{E} . It was found that these contributions to the longitudinal cross section σ_L are not sufficient for describing the physical observables in the π^0 production at sufficiently low Q^2 [23]. The essential contributions from the transversity GPDs H_T , \tilde{E}_T are needed [24] to be consistent with the experiment. Within the handbag approach, the transversity GPDs together with the twist-3 meson wave function [24] contribute to the amplitudes of transversely polarized photons, which produce the transverse cross section σ_T , which is much larger with respect to the leading twist σ_L .

We discuss the handbag approach and the properties of meson production amplitudes in Section II. We show that the transversity GPD contributions, which have the twist-3 nature, lead to a large transverse cross section.

In beginning of Section III, we investigate the role of transversity GPDs in the cross sections of the π^0 leptonproduction at CLAS and COMPASS energies and show that our results are in good agreement with the experiment. Subsequently, we perform predictions for π^0 cross sections at EicC energies.

II. HANDBAG APPROACH. PROPERTIES OF MESON PRODUCTION AMPLITUDES

In the handbag approach, the meson photoproduction amplitude is factorized into a hard subprocess amplitude \mathcal{H} and GPDs F , which include information on the hadron structure at sufficiently high Q^2 . Note that for the leading twist amplitudes with longitudinally polarized photons the factorization was proved [2, 3]. In what follows, we consider the twist-3 contributions from transversity GPDs H_T and \tilde{E}_T as well. Factorization of these twist-3 amplitudes is an assumption. The process of the handbag approach is shown in Fig. 1.

The subprocess amplitude is computed employing the modified perturbative approach (MPA) [25]. The power k_\perp^2/Q^2 corrections are considered in the propagators of the hard subprocess \mathcal{H} together with the nonperturbative k_\perp -dependent meson wave function [26, 27]. The power corrections can be treated as an effective consideration of the higher twist contribution. The gluonic corrections are regarded in the form of the Sudakov factors. Resummation of the Sudakov factor can be done in the impact parameter space [25].

The unpolarized $ep \rightarrow e\pi^0 p$ cross section can be decomposed into a number of partial cross sections, which are observables of the process $\gamma^* p \rightarrow \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi} = \frac{1}{2\pi} \left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos 2\phi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi \frac{d\sigma_{LT}}{dt} \right). \quad (1)$$

The partial cross sections are expressed in terms of the $\gamma^* p \rightarrow \pi^0 p$ helicity amplitudes. When we omit the small $M_{0,-,+}$ amplitude, they can be written as follows

$$\begin{aligned} \frac{d\sigma_L}{dt} &= \frac{1}{\kappa} \left[|M_{0+,0+}|^2 + |M_{0-,0+}|^2 \right], \\ \frac{d\sigma_T}{dt} &= \frac{1}{2\kappa} \left(|M_{0-,++}|^2 + 2|M_{0+,++}|^2 \right), \\ \frac{d\sigma_{LT}}{dt} &= -\frac{1}{\sqrt{2}\kappa} \text{Re} \left[M_{0-,++}^* M_{0-,0+} \right], \\ \frac{d\sigma_{TT}}{dt} &= -\frac{1}{\kappa} |M_{0+,++}|^2. \end{aligned} \quad (2)$$

with

$$\kappa = 16\pi(W^2 - m^2) \sqrt{\Lambda(W^2, -Q^2, m^2)}. \quad (3)$$

Here, $\Lambda(x, y, z)$ is defined as $\Lambda(x, y, z) = (x^2 + y^2 + z^2) - 2xy - 2xz - 2yz$.

The amplitudes can be written as

$$\begin{aligned} M_{0-,0+} &= \frac{e_0}{Q} \frac{\sqrt{-t'}}{2m} \langle \tilde{E} \rangle, \\ M_{0+,0+} &= \sqrt{1-\xi^2} \frac{e_0}{Q} [\langle \tilde{H} \rangle - \frac{\xi^2}{1-\xi^2} \langle \tilde{E} \rangle], \\ M_{0-,++} &= \frac{e_0}{Q} \sqrt{1-\xi^2} \langle H_T \rangle, \\ M_{0+,++} &= -\frac{e_0}{Q} \frac{\sqrt{-t'}}{4m} \langle \tilde{E}_T \rangle, \end{aligned} \quad (4)$$

where $e_0 = \sqrt{4\pi\alpha}$ with $\alpha = \frac{1}{137}$ is the fine structure con-

stant.

$$\xi = \frac{x_B}{2 - x_B} \left(1 + \frac{m_p^2}{Q^2} \right), \quad t' = t - t_0, \quad t_0 = -\frac{4m^2\xi^2}{1 - \xi^2}. \quad (5)$$

x_B is the Bjorken variable with $x_B = Q^2/(W^2 + Q^2 - m^2)$. m is the proton mass and m_p is the meson mass.

At the leading-twist accuracy, the PM production is only sensitive to the polarized GPDs \tilde{H} and \tilde{E} , which contribute to the amplitudes of longitudinally polarized virtual photons [23]. The $\langle F \rangle$ in Eq. (4) with $F = \tilde{H}, \tilde{E}$ are the convolutions of the hard scattering amplitude $\mathcal{H}_{0\mu',0+}$ and GPDs F

$$\langle F \rangle = \int_{-1}^1 dx \mathcal{H}_{0\mu',0+} F(x, \xi, t). \quad (6)$$

The hard part is calculated employing the k -dependent wave function [26, 26], describing the longitudinally polarized mesons. The amplitude \mathcal{H} is represented as the contraction of the hard part M , which can be computed perturbatively, and the non-perturbative meson wave function ϕ_M , which can be found in Ref. [23]

$$\mathcal{H}_{\mu',\mu+} = \frac{2\pi\alpha_s(\mu_R)}{\sqrt{2N_c}} \int_0^1 d\tau \int \frac{d^2k_\perp}{16\pi^3} \phi_M(\tau, k_\perp^2) M_{\mu'\mu}. \quad (7)$$

The GPDs are constructed adopting the double distribution representation [19]

$$F(x, \xi, t) = \int_{-1}^1 d\rho \int_{-1+|\rho|}^{1-|\rho|} d\gamma \delta(\rho + \xi\gamma - x) \omega(\rho, \gamma, t), \quad (8)$$

which connects GPDs F with PDFs h via the double distribution function ω . For the valence quark double distribution, it is

$$\omega(\rho, \gamma, t) = h(\rho, t) \frac{3}{4} \frac{[(1-|\rho|)^2 - \gamma^2]}{(1-|\rho|)^3}. \quad (9)$$

The t -dependence in PDFs h is presented in the Regge form

$$h(\rho, t) = N e^{(b-\alpha' \ln \rho)t} \rho^{-\alpha(0)} (1-\rho)^\beta, \quad (10)$$

and $\alpha(t) = \alpha(0) + \alpha' t$ is the corresponding Regge trajectory. The parameters in Eq. (10) are fitted from the known information about PDFs [28] e.g., or from the nucleon form factor analysis [29]. We consider the Q^2 evolution of GPDs through the evolution of the gluon distribution, as in Eq. (9); see [20]. The evolution was tested for valence quarks as well. It is approximately calculated for

the kinematical range in this work. We are working in the range $2 < Q^2 < 7 \text{ GeV}^2$. The parameters of GPDs are determined at the middle point $Q^2 = 4 \text{ GeV}^2$. In these very limited Q^2 range, the explicit form of the GPDs evolution is not so essential.

It was found that at low Q^2 , data on the PM lepton production also requires the contributions from the transversity GPDs H_T and $\bar{E}_T = 2\tilde{H}_T + E_T$, which determine the amplitudes $M_{0-,++}$ and $M_{0+,++}$, respectively. Within the handbag approach, the transversity GPDs are accompanied by a twist-3 meson wave function in the hard amplitude \mathcal{H} [24] which is the same for both the $M_{0\pm,++}$ amplitudes in Eq. (4). For the corresponding transversity convolutions, we have forms similar to (6), as follows:

$$\begin{aligned} \langle H_T \rangle &= \int_{-1}^1 dx \mathcal{H}_{0-,++}(x, \dots) H_T; \\ \langle \bar{E}_T \rangle &= \int_{-1}^1 dx \mathcal{H}_{0-,++}(x, \dots) \bar{E}_T. \end{aligned} \quad (11)$$

There is a parameter μ_p in the twist-3 meson wave function which is large and enhanced by the chiral condensate. In our calculation, we adopt $\mu_p = 2 \text{ GeV}$ at a scale of 2 GeV.

The H_T GPDs are connected with the transversity PDFs as follows:

$$h_T(\rho, 0) = \delta(\rho); \quad \text{and} \quad \delta(\rho) = N_T \rho^{1/2} (1-\rho) [q(\rho) + \Delta q(\rho)], \quad (12)$$

by employing the model [30]. We define the t -dependence of h_T as in Eq. (10).

The information on \bar{E}_T can be obtained now only in the lattice QCD [31]. The lower moments of \bar{E}_T^u and \bar{E}_T^d were found to be quite large, have the same sign and be of a similar size. As a result, we have large \bar{E}_T contributions to the π^0 production. This is parameterized by the form as in Eq. (10).

III. TRANSVERSITY EFFECTS IN π^0 MESON LEPTOPRODUCTION

In this section, we present our results on π^0 lepton production based on the handbag approach. In the calculation, we adopt the leading contribution Eq. (2) together with the transversity effects described in Eq. (11), which are essential at low Q^2 . The amplitudes are calculated based on the PARTONS collaboration code [32] that was modified in Fortran employing results of the GK model for GPDs [24].

In Fig. 2, we present the model results for the π^0 production cross section compared with the CLAS experimental data [13]. The transverse cross section, in which

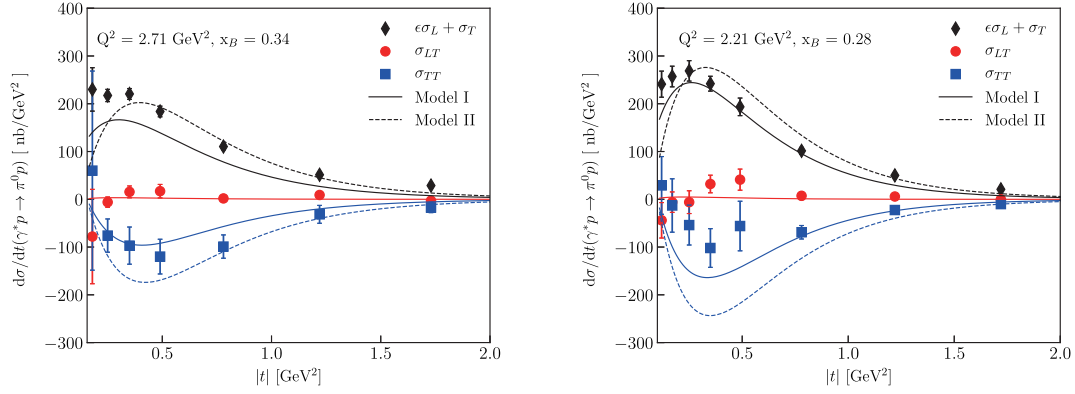


Fig. 2. (color online) Cross section of π^0 production in the CLAS energy range together with the data [13]. Black lines describe $\sigma = \sigma_T + \epsilon\sigma_L$, red lines represent σ_{LT} , blue lines depict σ_{TT} .

the \bar{E}_T and H_T contributions are important [24] dominates at low Q^2 . At small momentum transfer, the H_T effects are visible and provide a nonzero cross section. At $|t'| \sim 0.3 \text{ GeV}^2$, the \bar{E}_T contribution becomes essential in σ_T and results in a maximum for the cross section. A similar contribution from \bar{E}_T is observed in the interference cross section σ_{TT} [24]. For the calculations we use the parameters in Table 1. Details for \bar{H} parameterization can be found in [24]. The fact that we describe well both unseparated $\sigma = \sigma_T + \epsilon\sigma_L$ and σ_{TT} cross sections indicates that the transversity H_T and \bar{E}_T effects were observed at CLAS [13]. Note that in this experiment it was not possible to separate σ_L and σ_T . The model produces at CLAS kinematics the leading twist $\frac{d\sigma_L}{dt}(|t| = 0.3 \text{ GeV}^2) \sim \text{few nb/GeV}^2$. This is about two orders of magnitude smaller with respect to σ . Thus, we see that σ_T determined by the twist 3 effects gives a dominant contribution to unseparated σ . This prediction of the model [24] was confirmed by the JLab Hall A collaboration [33] by using the Rosenbluth separation of the π^0 electroproduction cross section.

Our results for COMPASS kinematics are shown in Fig. 3. It can be seen that Model I gives results about two times larger with respect to the COMPASS data [14]. That was the reason for changing to model parameters that permit describing both CLAS and COMPASS data. New parameters for Model II are exhibited in Table 2 [34]. Because \bar{E}_T contribution is essential in the σ_T and σ_{TT} cross sections, parameterization changes mainly the energy dependence of this GPD. Other GPDs are slightly changed to be consistent with experiments; see Fig. 2 and Fig. 3, in which both model results are shown.

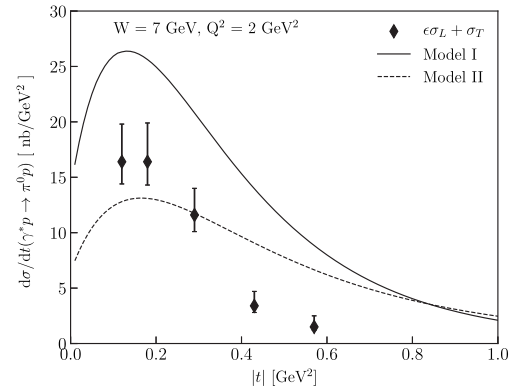


Fig. 3. Models results at COMPASS kinematics. Experimental data are from [14], solid curve is the prediction of Model I and dashed line presents the results of Model II.

The average COMPASS kinematics results for the cross sections are [14]

$$\begin{aligned} \langle \frac{d\sigma_{TT}}{dt} \rangle &= -(6.1 \pm 1.3 \pm 0.7) \text{ nb/GeV}^2, \\ \langle \frac{d\sigma_{LT}}{dt} \rangle &= (1.5 \pm 0.5 \pm 0.3) \text{ nb/GeV}^2. \end{aligned} \quad (13)$$

Model II gives the following results at the same kinematics

$$\begin{aligned} \langle \frac{d\sigma_{TT}}{dt} \rangle &= -6.4 \text{ nb/GeV}^2, \\ \langle \frac{d\sigma_{LT}}{dt} \rangle &= 0.1 \text{ nb/GeV}^2, \end{aligned} \quad (14)$$

Table 1. Regge parameters and normalizations of the GPDs at a scale of 2 GeV. Model I.

GPD	$\alpha(0)$	β^u	β^d	α'/GeV^{-2}	b/GeV^{-2}	N^u	N^d
\bar{E}	0.48	5	5	0.45	0.9	14.0	4.0
\bar{E}_T	0.3	4	5	0.45	0.5	6.83	5.05
H_T	—	—	—	0.45	0.3	1.1	-0.3

Table 2. Regge parameters and normalizations of the GPDs at a scale of 2 GeV. Model II.

GPD	$\alpha(0)$	α'/GeV^{-2}	b/GeV^{-2}	N^u	N^d
$\tilde{E}_{n.p.}$	0.32	0.45	0.6	18.2	5.2
\tilde{E}_T	-0.1	0.45	0.67	29.23	21.61
H_T	-	0.45	0.04	0.68	-0.186

which is close to the COMPASS results in Eq. (13). Model I gives cross sections that are about two times larger with respect to Model II. This is the same effect as we see in Fig. 3. This means that COMPASS provides essential constraints on the \tilde{E}_T contribution.

Using new GPDs parameterization may be important at EicC because its energy range lies not far from that of COMPASS. In future analyzes, we will give predictions for both GPDs models I and II, as at higher energies, a detailed study of transversity GPDs can be done.

In Fig. 4 and Fig. 5 we show the W and Q^2 dependencies of the σ and σ_{TT} cross sections in the EicC energy range. We show results for $W = 8, 12, 16$ GeV and $Q^2 = 2, 5, 7$ GeV² that are typical for EicC kinematics. The cross sections σ_{LT} are rather small and difficult to distinguish on these figures. Thus we separate them into individual Fig. 6 and Fig. 7, where the W and Q^2 dependencies of σ_{LT} are shown in pb/GeV². We use the same W and Q^2 values as for Fig. 4 and Fig. 5. One can see that

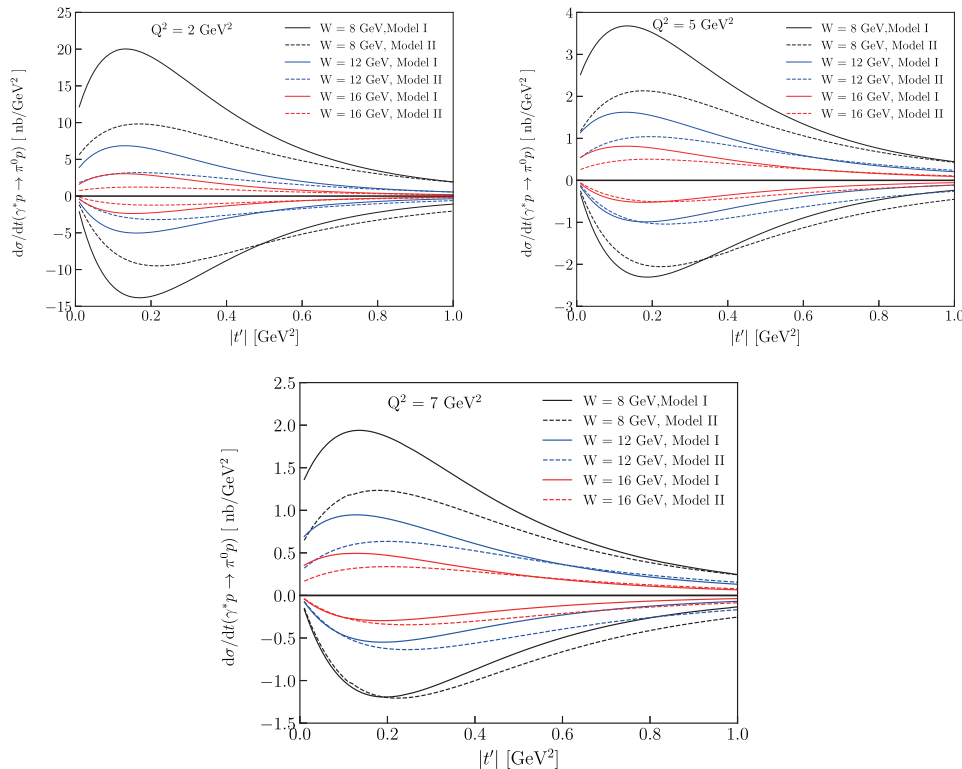
all cross sections decrease with increasing W and Q^2 . Model II gives typically smaller results with respect to Model I. At EicC kinematics, we get a rather small leading twist cross section σ_L , which is about one order of magnitude smaller with respect to σ_T . This means that the dominance of twist-3 transversity effects observed at low energy [24, 33] is valid up to high EicC energies. Our predictions on π^0 production give the possibility of performing a more detailed test on the energy dependencies of the transversity GPDs in future EicC experiments.

Now, we shall briefly discuss whether it is really possible to analyze the energy dependencies of transversity GPDs H_T and \tilde{E}_T from experimental data on cross sections. In experiments (see, e.g., [13]) usually, the unseparated cross sections $\sigma = \epsilon\sigma_L + \sigma_T$, σ_{LT} and σ_{TT} are measured. σ_L is determined by the twist-2 contribution. It is rather small and can be omitted in our estimations. Thus $\sigma \propto \sigma_T$ here. We will not discuss here σ_{LT} here.

We see that if

$$\frac{d\sigma_T}{dt} \sim -\frac{d\sigma_{TT}}{dt},$$

this means that in this range the essential contribution comes from the M_{0+++} amplitude (see (2)). At CLAS and COMPASS energies, this approximately happened at $|t'| = 0.3$ GeV². This means that at this momentum trans-

**Fig. 4.** (color online) Models results for $\sigma = \sigma_T + \epsilon\sigma_L$ and σ_{TT} cross section at EicC kinematics. W dependencies at fixed Q^2 are shown. The curves above the X -axis are predictions of σ , and the curves below the X -axis are predictions of σ_{TT} .

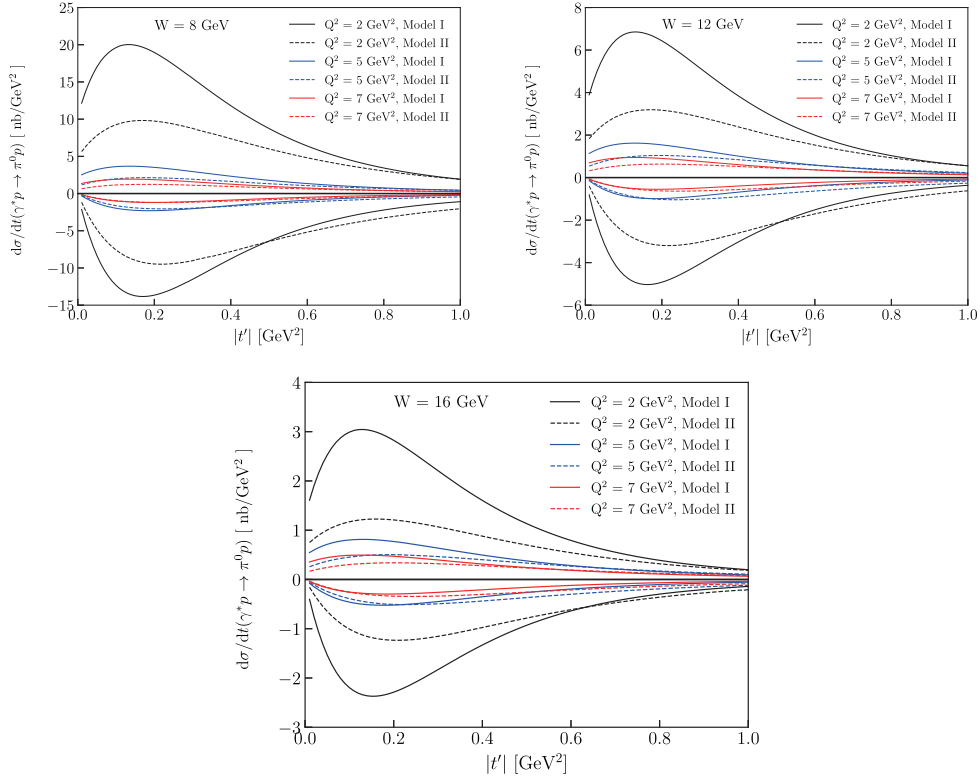


Fig. 5. (color online) Models results for $\sigma = \sigma_T + \epsilon\sigma_L$ and σ_{TT} cross sections at EicC kinematics. Q^2 dependencies at fixed W are shown. The curves above the X -axis are predictions of σ , and the curves below the X -axis are predictions of σ_{TT} .

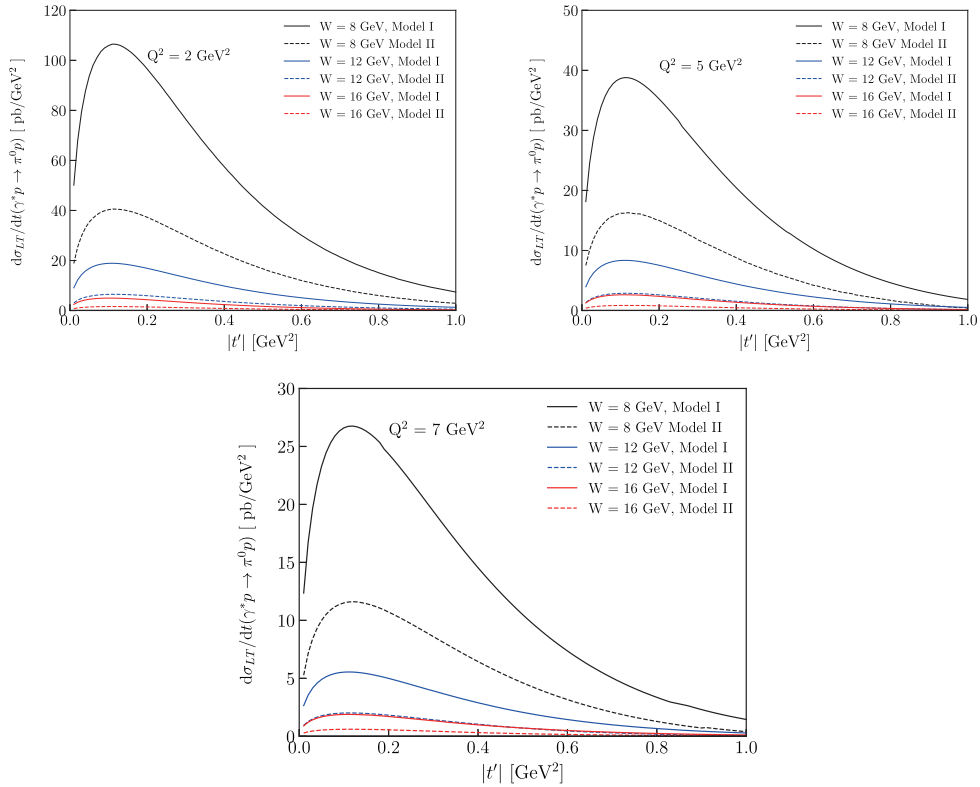


Fig. 6. (color online) Models predictions for σ_{LT} cross sections (in pb/GeV²) at EicC kinematics as a function of W at fixed Q^2 .

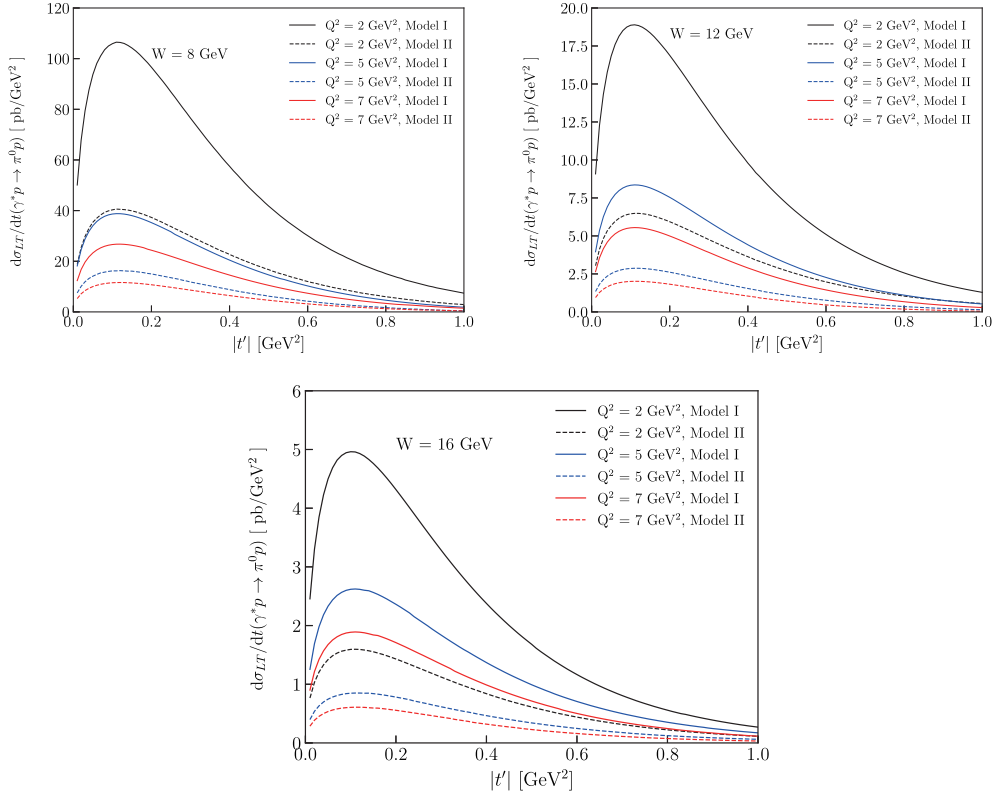


Fig. 7. (color online) Models results for σ_{LT} cross sections (in pb/GeV²) at EicC kinematics as a function of Q^2 at fixed W .

for the $\langle \bar{E}_T \rangle$ contribution dominates. At $|t'| = 0$ GeV² the \bar{E}_T is equal to zero. This means that at this point the $\langle H_T \rangle$ contribution is essential.

Thus, using Eqs. (2)–(4), we can determine two quantities

$$\begin{aligned} \langle H_T \rangle &\propto \sqrt{\kappa \frac{d\sigma_T}{dt}(|t'| = 0 \text{ GeV}^2)}, \\ \langle \bar{E}_T \rangle &\propto \sqrt{\kappa \frac{d\sigma_T}{dt}(|t'| = 0.3 \text{ GeV}^2)}, \end{aligned} \quad (15)$$

and one more in addition

$$\langle \bar{E}_T(TT) \rangle \propto \sqrt{\kappa \left| \frac{d\sigma_{TT}}{dt}(|t'| = 0.3 \text{ GeV}^2) \right|}. \quad (16)$$

Eq. (15) is a some approximation based on \bar{E}_T dominance near $|t'| \sim 0.3$ GeV². Eq. (16) gives direct information on \bar{E}_T , but $\frac{d\sigma_{TT}}{dt}$ is more difficult to study.

Thus, one can try to analyze the W dependencies of the cross sections at $|t'| \sim 0$ GeV² and $|t'| \sim 0.3$ GeV² to determine the energy dependencies of H_T and \bar{E}_T .

The results of the model calculations for the quantities in Eq. (15) for the GPDs models I and II can be parameterized as follows:

meterized as follows:

$$\langle H \rangle \sim A W^n. \quad (17)$$

We shall estimate n power using the results from Eq. (15) and n_H directly from the energy dependencies of the GPDs in the $W = 3 \sim 15$ GeV interval. The results are

$$\langle \bar{E}_T^{\text{Model-II}} \rangle: \quad n = 0.53, \quad n_H = 0.5; \quad (18)$$

$$\langle \bar{E}_T^{\text{Model-I}} \rangle: \quad n = 0.72, \quad n_H = 0.7; \quad (19)$$

$$\langle H_T \rangle: \quad n = 0.8, \quad n_H = 0.75. \quad (20)$$

We see that the energy dependencies of models II and I are rather different. From (16) we find the same power as in Eq. (19).

Thus, we find very closed powers n from cross section analyzes and directly from GPDs. This mean that we can really estimate the energy (s_B) dependencies of the GPDs from experimental data.

IV. CONCLUSION

The exclusive electroproduction of π^0 mesons was

analyzed here within the handbag approach, in which the amplitude was factorized into two parts. The first one are the subprocess amplitudes, which are calculated using k_{\perp} factorization [25]. The other essential ingredients are the GPDs, which contain information about the hadron structure. The results on the cross sections based on this approach were found to be in good agreement with data at HERMES, COMPASS energies at high Q^2 [24].

The leading-twist accuracy is not sufficient to describe π^0 leptonproduction at low Q^2 . It was confirmed [24] that rather strong transversity twist-3 contributions are required by the experiment. In the handbag approach, they are determined by the transversity GPDs H_T and \bar{E}_T in convolution with a twist-3 pion wave function. The transversity GPDs lead to a large transverse cross section for π^0 production.

Here, we consider two GPDs parameterizations. Model I was proposed in [24] to obtain a good description of the CLAS collaboration [13]. Later on the COMPASS ex-

periment produced π^0 data at higher energies [14]. Model I predictions at COMPASS energies are higher with respect to the experiment by a factor of 2. The energy dependencies of the transversity GPDs were modified in Model II [34], which describes properly both CLAS and COMPASS data.

In this analysis we perform predictions for unseparated σ , σ_{LT} and σ_{TT} cross sections for EicC kinematics for both models I and II. We confirm that the transversity dominance $\sigma_T \gg \sigma_L$, observed at low CLAS energies is valid up to the EicC energies range. Our results can be applied in future EicC experiments to give additional essential constraints on the transversity GPDs in the EicC energy range.

ACKNOWLEDGMENT

S. G. expresses his gratitude to P. Kroll for long-time collaboration on GPD studies.

References

- [1] D. Müller, D. Robaschik, B. Geyer *et al.*, *Fortsch. Phys.* **42**, 101-141 (1994), arXiv:hep-ph/9812448[hep-ph]
- [2] X. D. Ji, *Phys. Rev. Lett.* **78**, 610-613 (1997), arXiv:hep-ph/9603249[hep-ph]
- [3] A. V. Radyushkin, *Phys. Lett. B* **385**, 333-342 (1996), arXiv:hep-ph/9605431[hep-ph]
- [4] X. D. Ji, *Phys. Rev. D* **55**, 7114-7125 (1997), arXiv:hep-ph/9609381[hep-ph]
- [5] A. V. Radyushkin, *Phys. Rev. D* **56**, 5524-5557 (1997), arXiv:hep-ph/9704207[hep-ph]
- [6] J. J. C. Collins, L. Frankfurt and M. Strikman, *Phys. Rev. D* **56**, 2982-3006 (1997), arXiv:hep-ph/9611433[hep-ph]
- [7] K. Goeke, M. V. Polyakov, and M. Vanderhaeghen, *Prog. Part. Nucl. Phys.* **47**, 401-515 (2001), arXiv:hep-ph/0106012[hep-ph]
- [8] M. Vanderhaeghen, P. A. M. Guichon, and M. Guidal, *Phys. Rev. D* **60**, 094017 (1999), arXiv:hep-ph/9905372[hep-ph]
- [9] G. Duplancić, D. Müller and K. Passek-Kumerički, *Phys. Lett. B* **771**, 603-610 (2017), arXiv:1612.01937[hep-ph]
- [10] M. Constantinou, A. Courtoy, M. A. Ebert *et al.*, *Prog. Part. Nucl. Phys.* **121**, 103908 (2021), arXiv:2006.08636[hep-ph]
- [11] M. Diehl, *Phys. Rept.* **388**, 41-277 (2003), arXiv:hep-ph/0307382[hep-ph]
- [12] A. V. Belitsky and A. V. Radyushkin, *Phys. Rept.* **418**, 1-387 (2005), arXiv:hep-ph/0504030[hep-ph]
- [13] I. Bedlinskiy *et al.* (CLAS Collaboration), *Phys. Rev. C* **90**(2), 025205 (2014), arXiv:1405.0988[nucl-ex]
- [14] M. G. Alexeev *et al.* (COMPASS Collaboration), *Phys. Lett. B* **805**, 135454 (2020), arXiv:1903.12030[hep-ex]
- [15] A. Accardi, J. L. Albacete, M. Anselmino *et al.* *Eur. Phys. J. A* **52**(9), 268 (2016) arXiv: 1212.1701 [nucl-ex]
- [16] R. Abdul Khalek, A. Accardi, J. Adam *et al.*, *Nucl. Phys. A* **1026**, 122447 (2022)
- [17] D. P. Anderle, V. Bertone, X. Cao *et al.*, *Front. Phys.* (Beijing) **16**(6), 64701 (2021), arXiv:2102.09222[nucl-ex]
- [18] J. M. M. Chávez, V. Bertone, F. De Soto Borrero *et al.*, *Phys. Rev. Lett.* **128**(20), 202501 (2022), arXiv:2110.09462[hep-ph]
- [19] I. V. Musatov and A. V. Radyushkin, *Phys. Rev. D* **61**, 074027 (2000), arXiv:hep-ph/9905376[hep-ph]
- [20] S. V. Goloskokov and P. Kroll, *Eur. Phys. J. C* **42**, 281-301 (2005) arXiv: hep-ph/0501242
- [21] S.V. Goloskokov and P. Kroll, *Eur. Phys. J. C* **50**, 829-842 (2007)
- [22] S. V. Goloskokov and P. Kroll, *Eur. Phys. J. C* **59**, 809-819 (2009)
- [23] S. V. Goloskokov and P. Kroll, *Eur. Phys. J. C* **65**, 137-151 (2010), arXiv:0906.0460[hep-ph]
- [24] S. V. Goloskokov and P. Kroll, *Eur. Phys. J. A* **47**, 112 (2011), arXiv:1106.4897[hep-ph]
- [25] J. Botts and G. F. Sterman, *Nucl. Phys. B* **325**, 62-100 (1989)
- [26] R. Jakob and P. Kroll, *Phys. Lett. B* **315**, 463-470 (1993) [Erratum: *Phys. Lett. B* **319**, 545 (1993)] arXiv: hep-ph/9306259 [hep-ph]
- [27] J. Bolz, P. Kroll, and J. G. Körner, *Z. Phys. A* **350**, 145-159 (1994) arXiv: hep-ph/9403319 [hep-ph]
- [28] J. Pumplin, D. R. Stump, J. Huston *et al.*, *JHEP* **07**, 012 (2002), arXiv:hep-ph/0201195[hep-ph]
- [29] M. Diehl, T. Feldmann, R. Jakob *et al.*, *Eur. Phys. J. C* **39**, 1-39 (2005), arXiv:hep-ph/0408173[hep-ph]
- [30] M. Anselmino, M. Boglione, U. D'Alesio *et al.*, *Nucl. Phys. B Proc. Suppl.* **191**, 98-107 (2009), arXiv:0812.4366[hep-ph]
- [31] M. Göckeler *et al.* (QCDSF and UKQCD), *Phys. Rev. Lett.* **98**, 222001 (2007), arXiv:hep-lat/0612032[hep-lat]
- [32] B. Berthou, D. Binosi, N. Chouika *et al.*, *Eur. Phys. J. C* **78**(6), 478 (2018), arXiv:1512.06174[hep-ph]
- [33] M. Defurne *et al.* (Jefferson Lab Hall A), *Phys. Rev. Lett.* **117**(26), 262001 (2016), arXiv:1608.01003[hep-ex]
- [34] P. Kroll, private communication, (2019)