



Searching for dark matter with $t\bar{t}$ resonance

Yoav Afik^a, Eitan Gozani^b, Yoram Rozen^c

Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

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Abstract Many models containing particles which are candidates for dark matter, assume the standard model particles and the dark matter candidates are mediated by a spin-0 particle. At the LHC, one can use these models for dark matter searches. One of the possible approaches for the search of these models is by considering the decay of the spin-0 particle to a pair of $t\bar{t}$, thus modifying the pattern of the top quark pair invariant mass spectrum. This search suggests a good sensitivity in a parameter space different than the more traditional searches. We examine this sensitivity and put limits on two benchmark models containing candidates for dark matter, using previous ATLAS results. It was found that when the mediator mass (m_{Y_0}) and the dark matter candidate mass (m_χ) have values of $m_{Y_0} \sim 2 \cdot m_\chi$, mediator masses in the range of [400, 600] GeV are excluded. We compare our results to direct detection experiments and show that we gain sensitivity for new regions which are not covered by other searches.

1 Introduction

Astrophysical observations support the existence of nonbaryonic component of the universe: Dark Matter (DM) [1–4]. DM particles have to be stable, massive, and do not participate in the strong and electro-magnetic interactions. There are many searches for DM candidates at the Large Hadron Collider (LHC) experiment that use different approaches to model the signal for DM. One of the most popular candidates is a Weakly Interacting Massive Particle (WIMP) [5]. At the LHC, one can search for WIMP type DM particles (χ) produced in pp collisions.

Searches for DM using models with DM candidates and spin-0 mediators as a signal were already presented by

ATLAS [6–11] and CMS [12–16] collaborations, with up to $36 fb^{-1}$ of integrated luminosity, with centre of mass energy $\sqrt{s} = 13$ TeV. These searches focus on production of DM in association with a pair of top or bottoms quarks. In all of these searches, the mediator decays to a pair of DM candidates, leaving a signature of high missing transverse momentum in the detector.

When searching for dark matter, most of the analyses use some model as a benchmark. Although the conclusions of a given analysis are broader than the benchmark model used, in most of the cases, the analysis aims to maximise the sensitivity for this specific model. As general as the model is, parameter spaces of other models containing DM candidates are not necessarily covered. A complementary search for these models can be achieved if the mediator decays to a pair of top quarks, leaving a more complex signature in the detector. These searches are challenging as strong interference with the Standard Model (SM) $t\bar{t}$ production is expected [17], leading to a pick-dip shape in the spectrum of the $t\bar{t}$ invariant mass [18, 19]. Good sensitivity is expected if the mediator is heavier than twice the mass of the top quark. Some representative Feynman diagrams for leading-order production of a $t\bar{t}$ pair by a spin-0 mediator (Y_0) and by the SM are presented in Fig. 1.

There have been a few analyses targeting searches for heavy particles decaying to a pair of top quarks [20–22], creating a Breit-Wigner resonance in the $m_{t\bar{t}}$ spectrum. However, for most of those analyses, spin-0 particles were not taken into account. The recent search published in a similar

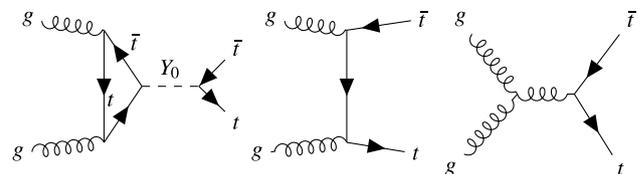


Fig. 1 Representative Feynman diagrams for a $t\bar{t}$ production via spin-0 mediator (Y_0) and via the SM

^a e-mail: yoavafik@campus.technion.ac.il

^b e-mail: gozani@campus.technion.ac.il

^c e-mail: rozen@physics.technion.ac.il

context, considering interference with the SM for $t\bar{t}$ production, was done by the ATLAS collaboration [23] using data of pp collisions at a centre-of-mass energy of $\sqrt{s} = 8$ TeV and integrated luminosity of $20.3 fb^{-1}$. This analysis used a Two-Higgs-Doublet model for the interpretation of the results, with a spin-0 particle mass of 500 – 750 GeV. Here, we show how the search for $t\bar{t}$ resonance originating from spin-0 particles is important for models containing DM. We cover parameter space that is not covered by the more traditional searches. Especially, searches for signature with a mediator that decays to a pair of DM candidates.

2 Theoretical framework

In many models containing new spin-0 particles, the couplings with the SM fermions are being set proportional to the SM Yukawa terms, by using the Minimal Flavour Violation (MFV) assumption [34]. This motivates a search in association with heavy flavour quarks. There are many models that assume interactions between DM candidates and spin-0 CP-odd or CP-even mediators, see for example [24, 27–33].

The width, at tree level, for a spin-0 particle decaying to a pair of Dirac fermions, which can be either DM candidates ($\chi\bar{\chi}$) or SM fermions ($f\bar{f}$), for models assuming MFV, is calculated as follows:

$$\Gamma(\phi/a \rightarrow \chi\bar{\chi}) = (g_{med-\chi\bar{\chi}})^2 \cdot \frac{m_{\phi/a}}{8\pi} \left(1 - \frac{4 \cdot m_{\chi}^2}{m_{\phi/a}^2}\right)^{n/2}, \tag{1}$$

$$\Gamma(\phi/a \rightarrow f\bar{f}) = (g_{med-f\bar{f}})^2 \cdot \frac{y_f^2 \cdot m_{\phi/a}}{16\pi} \left(1 - \frac{4 \cdot m_f^2}{m_{\phi/a}^2}\right)^{n/2}, \tag{2}$$

where $n = 3$ for a scalar (ϕ) and $n = 1$ for a pseudo-scalar (a). Here, χ is the DM candidate, f is the SM fermion, $m_{\phi/a}$ is the mass of the scalar / pseudo-scalar, m_{χ} is the DM mass, m_f and y_f are the corresponding mass and Yukawa term for the SM fermion, respectively. The parameters $g_{med-\chi\bar{\chi}}$ and $g_{med-f\bar{f}}$ are model dependent couplings. In general, Eq. 1 can be applied to other types of DM candidates, which are not Dirac fermions, but we keep it as a benchmark assumption. Interactions between the dark sector and the SM gauge bosons exist in part of those models, and are taken into account when analysing the results.

The calculation of the mediator width from Eqs. 1 and 2 presents an interesting behavior: If the mediator is heavy enough to decay to a pair of top quarks (i.e. $f\bar{f} = t\bar{t}$), the partial decay width of $\phi/a \rightarrow \chi\bar{\chi}$ becomes significantly smaller. This is especially true for high m_{χ} , where the partial decay width of $\phi \rightarrow \chi\bar{\chi}$ is highly suppressed. Therefore, the $t\bar{t}$ resonance search along the search for $t\bar{t} + \chi\bar{\chi}$

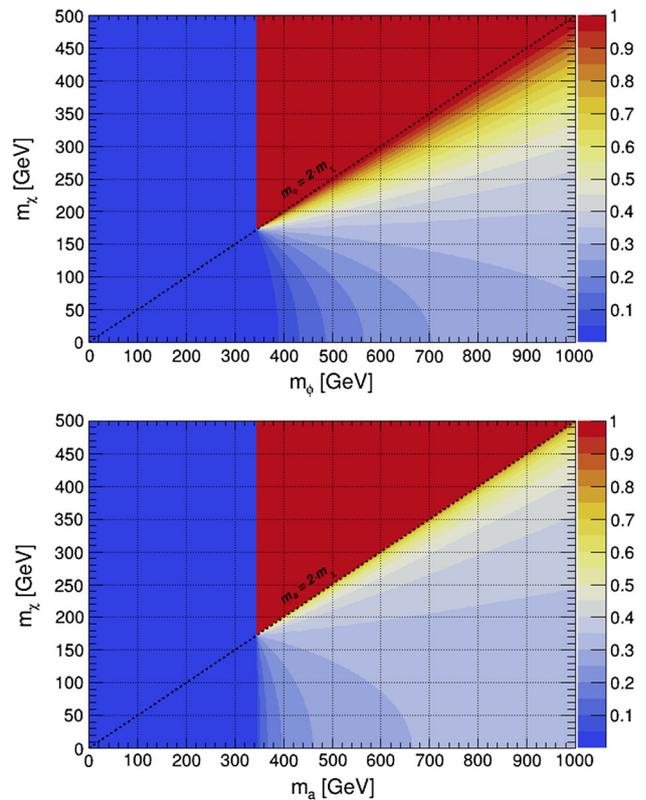


Fig. 2 Branching Ratio for the decay of a spin-0 mediator to a pair of $t\bar{t}$, assuming interactions with dark matter and tops only. The Branching Ratio is presented for a scalar (up) and pseudo-scalar (bottom) scenarios. The diagonal black dashed lines present the limit in the parameter space where the mediator is heavy enough to decay to a pair of DM candidates

are complementary: the former gains better sensitivity for low DM masses, while the latter has a better sensitivity for high DM masses. The Branching Ratio of this kind of spin-0 decay to a pair of top quarks is presented in Fig. 2, assuming couplings only for DM and top quarks, and setting $g_{med-\chi\bar{\chi}} = g_{med-f\bar{f}} = 1$. The behaviour discussed above is well observed in those figures.

3 Benchmark models

In order to emphasize how the behavior described in Sect. 2 affects the sensitivity for models containing a spin-0 mediator and DM candidates, we select two suitable benchmark models. The first model we consider is a simplified model with a spin-0 mediator that couples to both the SM fermions and to a new dark sector [24–26], where the couplings of the spin-0 mediator and the SM fermions are universal. The second model is a type-II two-Higgs-doublet model with an additional U(1) gauge symmetry, introduced in [27].

The first model describes a general set of couplings, and assumes that only fermions interact with the spin-0 media-

tors. The second model however, also describes interactions between the SM bosons and the spin-0 mediator, affecting the width of the spin-0 particle, and therefore modifying the invariant mass spectrum of the top quark pair. In addition, in the second model, unlike the first one, the couplings of the spin-0 mediator to the SM fermions are not universal, and therefore are different between down-type and up-type quarks (they are all fixed by one parameter, but they have different values).

3.1 Benchmark 1: simplified models

We consider two choices: in the first one the interaction with the SM is mediated by a scalar, while in the second scenario we consider only a new pseudo-scalar (assuming that the associated scalar is decoupled from the low-energy spectrum). The dark sector, in general, can contain more than a single particle. We assume the dark sector contain only one DM candidate which is a Dirac fermion. We keep in mind this assumption affects mostly the width of the mediator, so the results can be converted easily to more complicated cases. This model assumes Yukawa-like couplings between the dark sector mediator and the SM fermions. The interaction terms of the Lagrangians for the scalar (\mathcal{L}_ϕ) and pseudo-scalar (\mathcal{L}_a) are [24]:

$$\mathcal{L}_\phi^{int} = -g_\chi \phi \bar{\chi} \chi - \sum_{fermions} g_v \frac{y_f}{\sqrt{2}} \phi \bar{f} f, \tag{3}$$

$$\mathcal{L}_a^{int} = -i g_\chi a \bar{\chi} \gamma^5 \chi - \sum_{fermions} i g_v \frac{y_f}{\sqrt{2}} a \bar{f} \gamma^5 f. \tag{4}$$

Here, ϕ and a are scalar and pseudo-scalar fields which connect the SM with the dark sector, χ is the DM field, g_χ is the DM-mediator coupling, g_v is the flavour-universal SM-mediator coupling and y_f are the SM Yukawa couplings for fermions.

If $g_v \sim g_\chi$ and $m_{\phi/a} > 2m_\chi$, the decay of the mediator to DM is expected to dominate, unless the mediator is heavy enough for the top channel to open. This holds true because the Yukawa couplings to light fermions are significantly smaller compared to the Yukawa term of the top. The minimal viable value of $\Gamma_{\phi/a}$ can be calculated from the other parameters. The mediator width can be larger than the minimal one if additional dark sector particles are present. In our interpretation, however, we assume only one type of DM candidate. For simplicity, we use $g = g_\chi = g_v$, reducing the free parameters to three: $m_\chi, m_{\phi/a}, g$. To simulate the signal, the DMSIMP [35–37] models have been used with MadGraph [38].

3.2 Benchmark 2: 2HDM + Z'

This model is an extension to the familiar type-II 2HDM model [39], and introduces an extra spin-1 mass eigenstate which is denoted as Z' . The pseudo-scalar (A_0) is the only one that couples to a pair of DM candidates, therefore this particle is identified as the mediator between the SM particles and the dark sector. The decoupling limit [40] is assumed. Limits were set recently on this model by ATLAS [11], using final state with the SM Higgs boson (decays to a pair of bottom quarks or photons) and A_0 (decays to a pair of DM candidates). However, the parameter choice was set in order to emphasize the sensitivity of the signatures above. Therefore, we choose different parameters, in order to emphasize the final state discussed in this paper, and to add new constraints on the model.

This model has six parameters, which are set as follows: The mass of the light scalar, recognised as the SM Higgs boson in the decoupling limit, is set to $m_h = 125$ GeV; In order to avoid interference effects between two bosons, we choose the masses of the spin-0 bosons to allow only one of them, A_0 , to decay to a $t\bar{t}$ pair. Therefore, the mass of the heavier scalar and charged scalar are set equal to each other, $m_H = m_{H^\pm} = 300$ GeV, and the mass of the pseudo-scalar is set to $m_{A_0} = 400$ GeV; The parameters of the new Z' boson, its mass ($m_{Z'}$) and coupling ($g_{Z'}$), has a negligible effect on the $t\bar{t}$ invariant mass spectrum, and are chosen randomly to be $m_{Z'} = 3$ TeV and $g_{Z'} = 1$. The results are valid for other values of $m_{Z'}$, as long as it is heavy enough to avoid decays of A_0 to Z' with another Higgs boson; The ratio between vacuum expectation values of the two Higgs doublets, $\tan(\beta)$, and the mass of the DM candidate, m_χ , are free parameters. Those two parameters were selected to be scanned, since they have a direct impact on the strength of A_0 interactions and decay width, and therefore on the top pair invariant mass spectrum.

4 Results

Limits on spin-0 mediator models were already set at 95% Confidence Level (CL) [17], using the latest ATLAS $t\bar{t}$ resonance search with available data [21]. ATLAS also set limits on spin-0 mediators [23]. In both cases interference effects with the SM were considered, and the signal was modeled at Next to Leading Order (NLO) and Next to Next to Leading Order (NNLO) in QCD corrections, respectively. For those limits, however, no interaction with DM candidates was taken into account. This type of interaction is being considered in this section. The results of the former were found to be more efficient for mediator masses which are lower than 500 GeV, while the latter is more efficient for mediator masses which are higher than 500 GeV. This is expected since the corre-

sponding ATLAS analysis targeted spin-0 particle masses which are higher than 500 GeV.

The experimental resolution on the $t\bar{t}$ invariant mass, $m_{t\bar{t}}$, was calculated to be 8% for both analyses. Since the width of the mediator has a strong effect on the shape of the pure signal and interference distributions, an upper limit $\frac{\Gamma_{total}}{m_{\phi/a}} < 8\%$ was set, where Γ_{total} is the total decay width of the mediator. In the case that the mediator decays to a pair of DM candidates, $\frac{\Gamma_{total}}{m_{\phi/a}} < 40\%$ was used to keep the narrow width approximation valid. Results with higher total widths were discarded.

4.1 Benchmark 1: simplified models

Figure 3 presents upper limits at 95% CL on the coupling g . The figure presents the lowest coupling excluded for the model. Both scalar and pseudo-scalar mediator cases are considered. The best limits obtained from [10] searching for $t\bar{t} + \chi\bar{\chi}$ processes are presented for comparison. The exclusion contour is more stringent for the $t\bar{t}$ resonance searches when $m_{\phi/a} \geq 400$ GeV, especially when the DM mass is high. The limits obtained from the $t\bar{t}$ resonance are stronger for the scalar case since the width calculation (see Eqs. 1, 2) allows higher values for the pseudo-scalar case with similar parameters, leading to higher total widths which we discard.

The areas above the diagonal black dashed lines at Fig. 3 present the parameter space where the mediator is virtual for the production of $t\bar{t} + \chi\bar{\chi}$ processes ($m_{\phi/a} < 2 \cdot m_\chi$). In this part of the parameter space, comparing to the on-shell case (the area below the line with $m_{\phi/a} > 2 \cdot m_\chi$), the sensitivity for $t\bar{t} + \chi\bar{\chi}$ processes decreases significantly. This happens mainly because the cross section is lower (see for example Figure 6 at [10]).

The regions of the parameter space where the mediator mass is higher than twice the mass of the top ($m_{\phi/a} > 2 \cdot m_t$) are covered by the signature discussed in this paper for two cases: The first case is above the diagonal line, where the mass of the DM does not play a significant role, since in this case the mediator width is completely dominated by the top pair decay; the second case is below the diagonal line, allowing on-shell decay of the mediator both to a pair of top quarks and a pair of DM candidates. The limits in the second case are set close to the diagonal line, since the partial decay width of the mediator to $\chi\bar{\chi}$ decreases when $2 \cdot m_\chi$ increases up to $m_{\phi/a}$, as concluded from Eq. 1.

For both of the regions in the parameter space discussed in the previous paragraph, the $t\bar{t} + \chi\bar{\chi}$ search is not sensitive. However, the signature discussed in this paper is sensitive. On the contrary, in the case that the mediator is not heavy enough to decay to a pair of tops ($m_{\phi/a} < 2 \cdot m_t$) but decays to a pair of DM candidates ($m_{\phi/a} > 2 \cdot m_\chi$), the $t\bar{t} + \chi\bar{\chi}$ search is sensitive while the final state discussed in this paper is not.

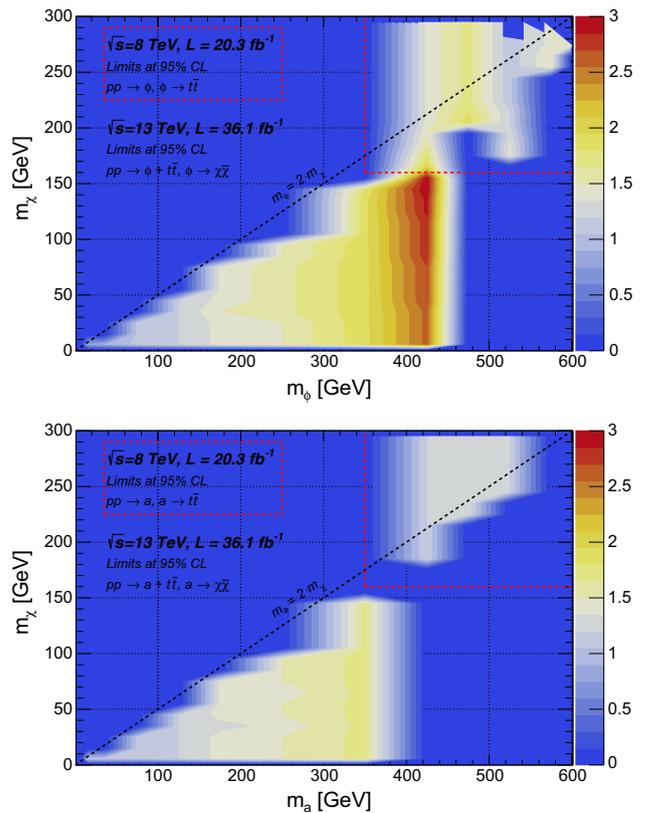


Fig. 3 Exclusion contour for DM simplified models in the $(m_{\phi/a}, m_\chi)$ plane for the scalar (upper) and pseudo-scalar (bottom) scenarios. The contour corresponds to the lowest value of the coupling $g = g_\chi = g_v$ allowed. Results from both $t\bar{t}$ resonance and $t\bar{t} + \chi\bar{\chi}$ signatures are presented for comparison, while the red dashed lines delimitate the limits set by the $t\bar{t}$ resonance signature. The diagonal black dashed lines present the limit in the parameter space where the mediator is heavy enough to decay to a pair of DM candidates

This emphasize the necessity of those signatures as being complementary to each other for this model.

4.2 Benchmark 2: 2HDM + Z'

Figure 4 presents the exclusion contour at 95% CL in the plane of $(m_\chi, \tan\beta)$. The remaining parameters of the model are fixed as discussed in Sect. 3. The best limits obtained from other searches are not presented since they do not provide any constraint in the considered parameter space. In this benchmark model the coupling between the top quarks and A_0 is proportional to $(\tan\beta)^{-1}$. Therefore, the lower part of the exclusion contour is due to mediator width values which are higher than 8% (see Eq. 2 with $g_{med-f\bar{f}} = (\tan\beta)^{-1}$), while the upper part is due to lower cross sections not excluded by the analysis.

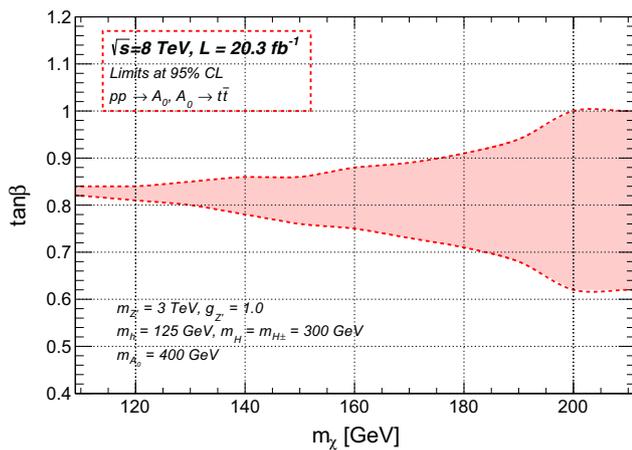


Fig. 4 Exclusion contour for 2HDM+Z' model in the $(m_\chi, \tan\beta)$ plane. The contour corresponds to the excluded values of the model. All of the other parameters of the model are fixed and mentioned in the figure. Only results using the $t\bar{t}$ resonance signature are presented, since other signatures do not exclude any part of the presented parameter space

4.3 Comparison to direct detection experiments

Results from LHC analyses with DM simplified models can be compared to direct detection experiments, for both spin-independent and spin-dependent cases. For this purpose, we use the prescription described at [41], setting limits on DM-nucleon interaction.

4.3.1 Spin-independent

For scalar simplified models, we set an upper limit on the DM-nucleon cross section (σ_{SI}) as follows [41]:

$$\sigma_{SI} \simeq 6.9 \cdot 10^{-43} \cdot \text{cm}^2 \cdot (g_\nu \cdot g_\chi)^2 \left(\frac{125 \text{ GeV}}{m_\phi}\right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}}\right)^2, \tag{5}$$

where we introduce the nucleon mass, $m_n = 0.939 \text{ GeV}$ (for both Neutrons and Protons), and the reduced DM-nucleon mass, $m_n \cdot m_\chi / (m_n + m_\chi)$. In the results presented at 4.1, we find that scalar masses in the range of $m_\phi \in [400, 600] \text{ GeV}$ are excluded, with the lowest couplings in the range of $g \in [1.1 - 2.0]$ and $m_\chi \geq 160 \text{ GeV}$. In order to use the prescription above, the couplings should be fixed. Therefore we choose $g_\chi = g_\nu = 1.5$. The results obtained in this paper are presented at Fig. 5, along with direct detection experiments [42–46] and with the contour calculated at [10]. There are two caveats for this comparison, however: the results we state are at 95% CL, while the results of the other experiments are at 90% CL; furthermore, one should keep in mind that the comparison is model dependent.

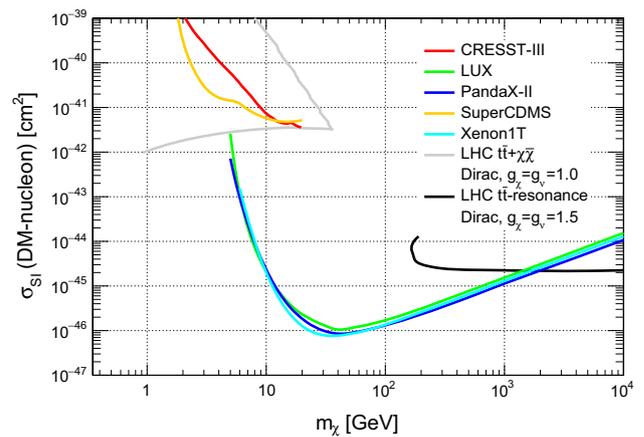


Fig. 5 Upper limit on spin-independent DM-nucleon cross section (σ_{SI}) as a function of the DM mass. The exclusion limit of the $t\bar{t}$ resonance obtained in this paper at 95% CL, is compared with limits from direct detection experiments and from the latest $t\bar{t} + \chi\bar{\chi}$ by ATLAS at 90% CL

4.3.2 Spin-dependent

For the pseudo-scalar scenario, a velocity suppression term in the non-relativistic limit creates large difference by several orders of magnitude in favour of LHC results [10,41], and therefore it is not presented. However, this actually makes the motivation of using $t\bar{t}$ resonance interpretation for DM models even stronger: it covers regions with higher DM masses, for which both $t\bar{t} + \chi\bar{\chi}$ and direct detection searches are insensitive to.

5 Conclusion

The necessity of $t\bar{t}$ resonance search as a complementary measurement for models containing DM candidates has been emphasized for two different models. Despite the challenges of this signature from the experimental point of view, it provides unique access to regions in the parameter space with high DM candidates masses, where other signatures quite often do not have the sensitivity. Further expansion of the parameter space covered by this signature is expected, using centre of mass energy of $\sqrt{s} = 13 \text{ TeV}$, and with the higher luminosity now recorded by ATLAS and CMS collaborations.

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References

1. F. Zwicky, Die Rotverschiebung von extragalaktischen Nebeln. *Helv. Phys. Acta* **6**, 110–127 (1933)
2. G. Bertone, D. Hooper, J. Silk, Particle dark matter: evidence, candidates and constraints. *Phys. Rept.* **405**, 279–390 (2005)
3. G. Jungman, M. Kamionkowski, K. Griest, Supersymmetric dark matter. *Phys. Rept.* **267**, 195–373 (1996)
4. J. Binney, S. Tremaine, Errata in Binney and Tremaine, 'Galactic Dynamics'. [arXiv:astro-ph/9304010](https://arxiv.org/abs/astro-ph/9304010) [astro-ph]
5. G. Steigman, M.S. Turner, Cosmological constraints on the properties of weakly interacting massive particles. *Nucl. Phys. B* **253**, 375 (1985)
6. ATLAS Collaboration, Search for a scalar partner of the top quark in the jets + E_T^{miss} Final State at $\sqrt{s} = 13$ TeV with the ATLAS detector. ATLAS-CONF-2016-077 (2016)
7. ATLAS Collaboration, Search for top squarks in final states with one isolated lepton, jets, and missing transverse momentum in $\sqrt{s} = 13$ TeV pp collisions with the ATLAS detector. ATLAS-CONF-2016-050 (2016)
8. ATLAS Collaboration, Search for direct top squark pair production and dark matter production in final states with two leptons in $\sqrt{s} = 13$ TeV pp collisions using 13.3 fb⁻¹ of ATLAS data. ATLAS-CONF-2016-076 (2016)
9. ATLAS Collaboration, Search for dark matter production associated with bottom quarks in 13.3 fb⁻¹ of pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector at the LHC. ATLAS-CONF-2016-086 (2016)
10. A. Morad, and others, Search for dark matter produced in association with bottom or top quarks in $\sqrt{s} = 13$ TeV pp collisions with the ATLAS detector. *Eur. Phys. J. C* **78**, 18 (2018)
11. ATLAS Collaboration, Constraints on mediator-based dark matter models using $\sqrt{s} = 13$ TeV pp collisions at the LHC with the ATLAS detector. ATLAS-CONF-2018-051 (2018)
12. A.M. Sirunyan, and others [CMS Collaboration], Search for dark matter produced in association with heavy-flavor quarks in proton-proton collisions at $\sqrt{s}=13$ TeV. *Eur. Phys. J. C* **77**, 845 (2017)
13. CMS Collaboration, Search for dark matter in association with a top quark pair at $\sqrt{s}=13$ TeV. CMS-PAS-EXO-16-005 (2016)
14. CMS Collaboration, Search for dark matter in association with a top quark pair at $\sqrt{s} = 13$ TeV in the dilepton channel. CMS-PAS-EXO-16-028 (2016)
15. CMS Collaboration, Search for Dark Matter produced in association with bottom quarks. CMS-PAS-B2G-15-007 (2016)
16. CMS Collaboration, Search for dark matter particles produced in association with a top quark pair at $\sqrt{s} = 13$ TeV. CMS-EXO-16-049 (2018)
17. B. Hespel, F. Maltoni, E. Vryonidou, Signal background interference effects in heavy scalar production and decay to a top-anti-top pair. *J. High Energy Phys.* **2016**, 16 (2016)
18. D. Dicus, A. Stange, S. Willenbrock, Higgs decay to top quarks at hadron colliders. *Phys. Lett. B* **333**, 126–131 (1994)
19. R. Frederix, F. Maltoni, Top pair invariant mass distribution: a window on new physics. *J. High Energy Phys.* **01**, 047 (2009)
20. CMS Collaboration, Search for anomalous $t\bar{t}$ production in the highly-boosted all-hadronic final state. *J. High Energy Phys.* **2012**, 29 (2012)
21. ATLAS Collaboration, A search for $t\bar{t}$ resonances using lepton-plus-jets events in proton-proton collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector. *J. High Energy Phys.* **2015**, 148 (2015)
22. ATLAS Collaboration, Search for heavy particles decaying into top-quark pairs using lepton-plus-jets events in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector. *Eur. Phys. J. C* **78**, 565 (2018)
23. ATLAS Collaboration, Search for heavy higgs bosons A/H decaying to a top quark pair in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector. *Phys. Rev. Lett.* **119**, 191803 (2015)
24. M.R. Buckley, D. Feld, D. Goncalves, Scalar simplified models for dark matter. *Phys. Rev. D* **91**, 015017 (2015)
25. U. Haisch, E. Re, Simplified dark matter top-quark interactions at the LHC. *J. High Energy Phys.* **2015**, 78 (2015)
26. A. Daniel, and others, Dark matter benchmark models for early LHC Run-2 searches: report of the ATLAS/CMS dark matter forum. [arXiv:1507.00966](https://arxiv.org/abs/1507.00966) [hep-ex]
27. A. Berlin, T. Lin, L.-T. Wang, Mono-Higgs detection of dark matter at the LHC. *J. High Energy Phys.* **2014**, 78 (2014)
28. M. Bauer, U. Haisch, F. Kahlhoefer, Simplified dark matter models with two Higgs doublets: I. Pseudoscalar mediators. *J. High Energy Phys.* **2017**, 138 (2017)
29. N.F. Bell, G. Busoni, I.W. Sanderson, Two Higgs doublet dark matter portal. *JCAP* **1801**(01), 015 (2018)
30. G. Arcadi, 2HDM portal for singlet-doublet dark matter. *Eur. Phys. J. C* **78**, 864 (2018)
31. L. Wang, X.-F. Han, B. Zhu, Light scalar dark matter extension of the type-II two-Higgs-doublet model. *Phys. Rev. D* **98**, 035024 (2018)
32. P. Bandyopadhyay, E.J. Chun, R. Mandal, Scalar Dark matter in leptophilic two-higgs-doublet model. *Phys. Lett. B* **779**, 201–205 (2018)
33. S. Dutta, A. Goyal, L.K. Saini, Spin-0[±] portal induced dark matter. *J. High Energy Phys.* **2018**, 23 (2018)
34. G. D'Ambrosio, G.F. Giudice, G. Isidori, A. Strumia, Minimal flavor violation: an effective field theory approach. *Nucl. Phys. B* **645**, 155–187 (2002)
35. O. Mattelaer, E. Vryonidou, Dark matter production through loop-induced processes at the LHC: the s-channel mediator case. *Eur. Phys. J. C* **75**, 436 (2015)
36. M. Backovic, M. Krämer, F. Maltoni, A. Martini, K. Mawatari, M. Pellen, Higher-order QCD predictions for dark matter production at the LHC in simplified models with s-channel mediators. *Eur. Phys. J. C* **75**, 482 (2015)
37. Y. Afik, F. Maltoni, K. Mawatari, P. Pani, G. Polesello, Y. Rozen, M. Zaro, DM+ $b\bar{b}$ simulations with DMSimp: an update. [arXiv:1811.08002](https://arxiv.org/abs/1811.08002) [hep-ex]
38. J. Alwall, M. Herquet, F. Maltoni, O. Mattelaer, T. Stelzer, MadGraph 5 : Going beyond. [arXiv:1106.0522](https://arxiv.org/abs/1106.0522) [hep-ph]
39. G.C. Branco, P.M. Ferreira, L. Lavoura, M.N. Rebelo, M. Sher, J.P. Silva, Theory and phenomenology of two-Higgs-doublet models. *Phys. Rept.* **516**, 1–102 (2012)
40. J.F. Gunion, H.E. Haber, The CP conserving two Higgs doublet model: the approach to the decoupling limit. *Phys. Rev. D* **67**, 075019 (2003)
41. A. Boveia, O. Buchmueller, C. Doglioni, K. Hahn, U. Haisch, F. Kahlhoefer, M. Mangano, C. McCabe, T.M.P. Tait, Recommendations on presenting LHC searches for missing transverse energy signals using simplified s-channel models of dark matter. [arXiv:1603.04156](https://arxiv.org/abs/1603.04156) [hep-ex]
42. D.S. Akerib, and others, Results from a search for dark matter in the complete LUX exposure. *Phys. Rev. Lett.* **118**, 021303 (2017)

43. Xiangxiang Ren, and others, Constraining Dark Matter Models with a Light Mediator at the PandaX-II Experiment, *Phys. Rev. Lett.* 121, 021304
44. E. Aprile, and others, First Dark Matter Search Results from the XENON1T Experiment, *Phys. Rev. Lett.* 119, 181301
45. R. Agnese, and others, New Results from the Search for Low-Mass Weakly Interacting Massive Particles with the CDMS Low Ionization Threshold Experiment, *Phys. Rev. Lett.* 116, 071301
46. F. Petricca, and others, First results on low-mass dark matter from the CRESST-III experiment, [arXiv:1711.07692](https://arxiv.org/abs/1711.07692) [astro-ph.CO]